

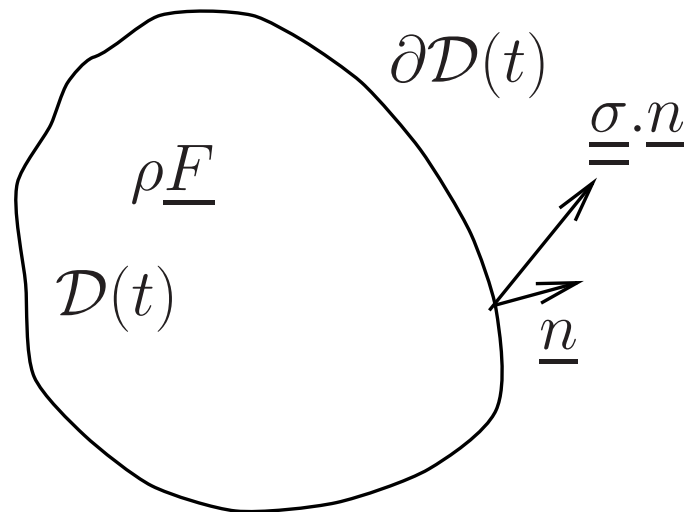
# FUNDAMENTALS OF FLUID MECHANICS

MMI103

Lesson 3 (course and exercises)

Thursday 25th September 2025

## MOMENTUM CONSERVATION LAW (reminder; Lesson 2)



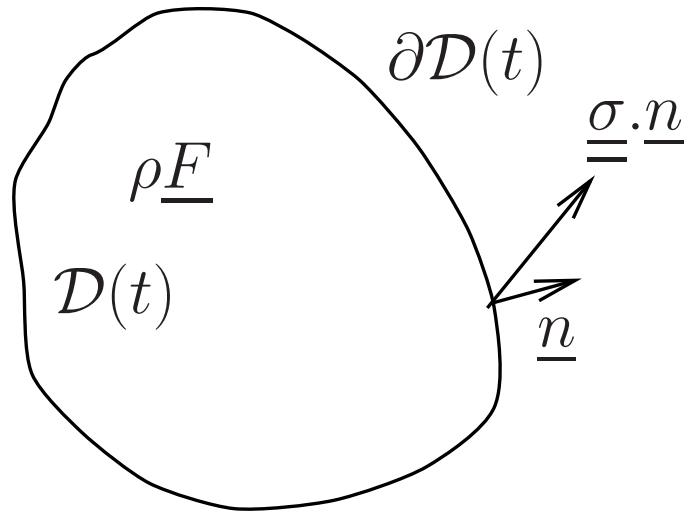
$$\frac{d}{dt} \int_{\mathcal{D}(t)} \rho \underline{u} d\Omega = \int_{\mathcal{D}(t)} \rho \underline{F} d\Omega + \int_{\partial\mathcal{D}(t)} \underline{\underline{\sigma}} \cdot \underline{n} da$$

$$\frac{\delta}{\delta t} \int_{\mathcal{D}} \rho \underline{u} d\Omega = \int_{\mathcal{D}} \rho \underline{F} d\Omega + \int_{\partial\mathcal{D}} [\underline{\underline{\sigma}} \cdot \underline{n} - \rho(\underline{u} \cdot \underline{n})\underline{u}] da$$

$\underline{\underline{\sigma}}$  symmetric,  $\rho \frac{d\underline{u}}{dt} = \rho \underline{F} + \text{div}(\underline{\underline{\sigma}})$ ,  $\text{div}(\underline{\underline{\sigma}}) = \sigma_{ij,j} \underline{e}_i$ ,  $\sigma_{ij,j} = \frac{\partial \sigma_{ij}}{\partial x_j}$

$$\frac{d\underline{u}}{dt} = \frac{\partial \underline{u}}{\partial t} + (u_{i,j} u_j) \underline{e}_i, \quad \frac{d\underline{u}}{dt} = \frac{\partial \underline{u}}{\partial t} + \underline{\underline{grad}}\left[\frac{u^2}{2}\right] + \underline{\underline{rot}}(\underline{u}) \wedge \underline{u}$$

# FLOWING NON-VISCOUS FLUID: EULER EQUATIONS



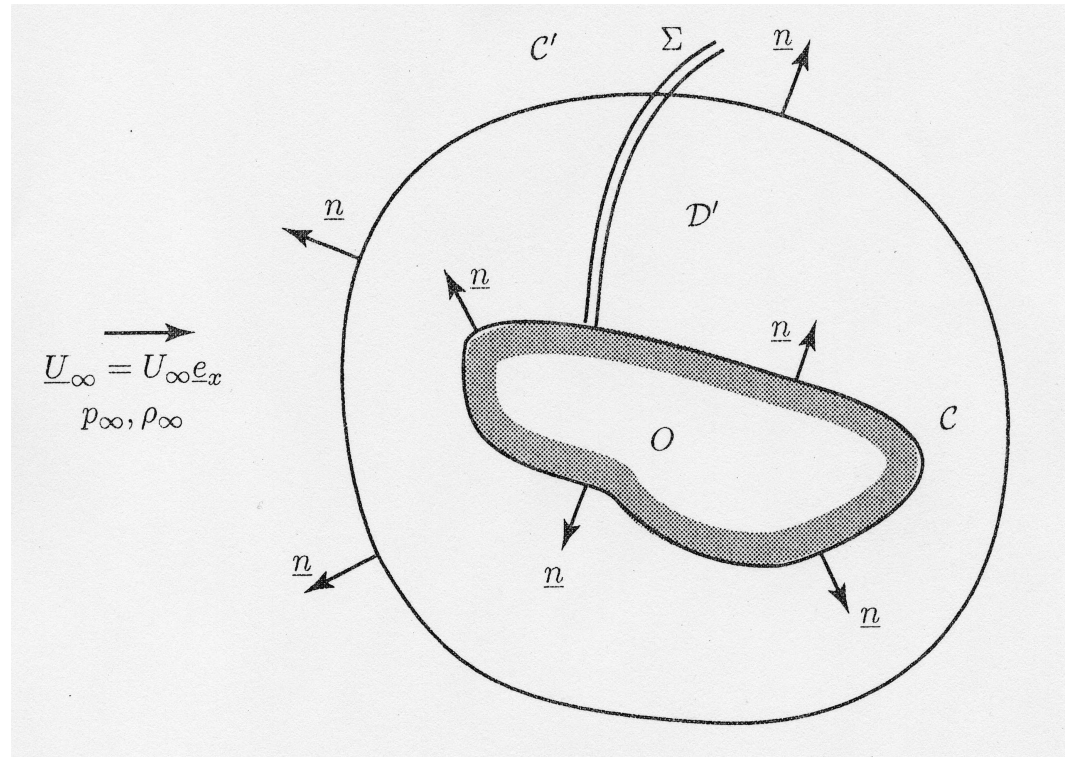
$$\frac{d}{dt} \int_{\mathcal{D}(t)} \rho \underline{u} d\Omega = \int_{\mathcal{D}(t)} \rho \underline{F} d\Omega - \int_{\partial\mathcal{D}(t)} p \underline{n} da$$

$$\frac{\delta}{\delta t} \int_{\mathcal{D}} \rho \underline{u} d\Omega = \int_{\mathcal{D}} \rho \underline{F} d\Omega - \int_{\partial\mathcal{D}} [p \underline{n} + \rho (\underline{u} \cdot \underline{n}) \underline{u}] da$$

$$\underline{\underline{\sigma}} = -p \underline{\underline{1}} \text{ (symmetric)}, \quad \frac{d\rho}{dt} + \rho \operatorname{div}(\underline{u}) = 0, \quad \rho \frac{d\underline{u}}{dt} = \rho \underline{F} - \underline{\underline{\operatorname{grad}}}[p],$$

Different cases: incompressible homogenous fluid:  $\rho = cste$ , compressible barotropic flow:  $\rho = \rho(p)$ ,  
 general  $\rho = \rho(p, T)$

# Force exerted by a steady flow of a non-viscous fluid on a motionless body

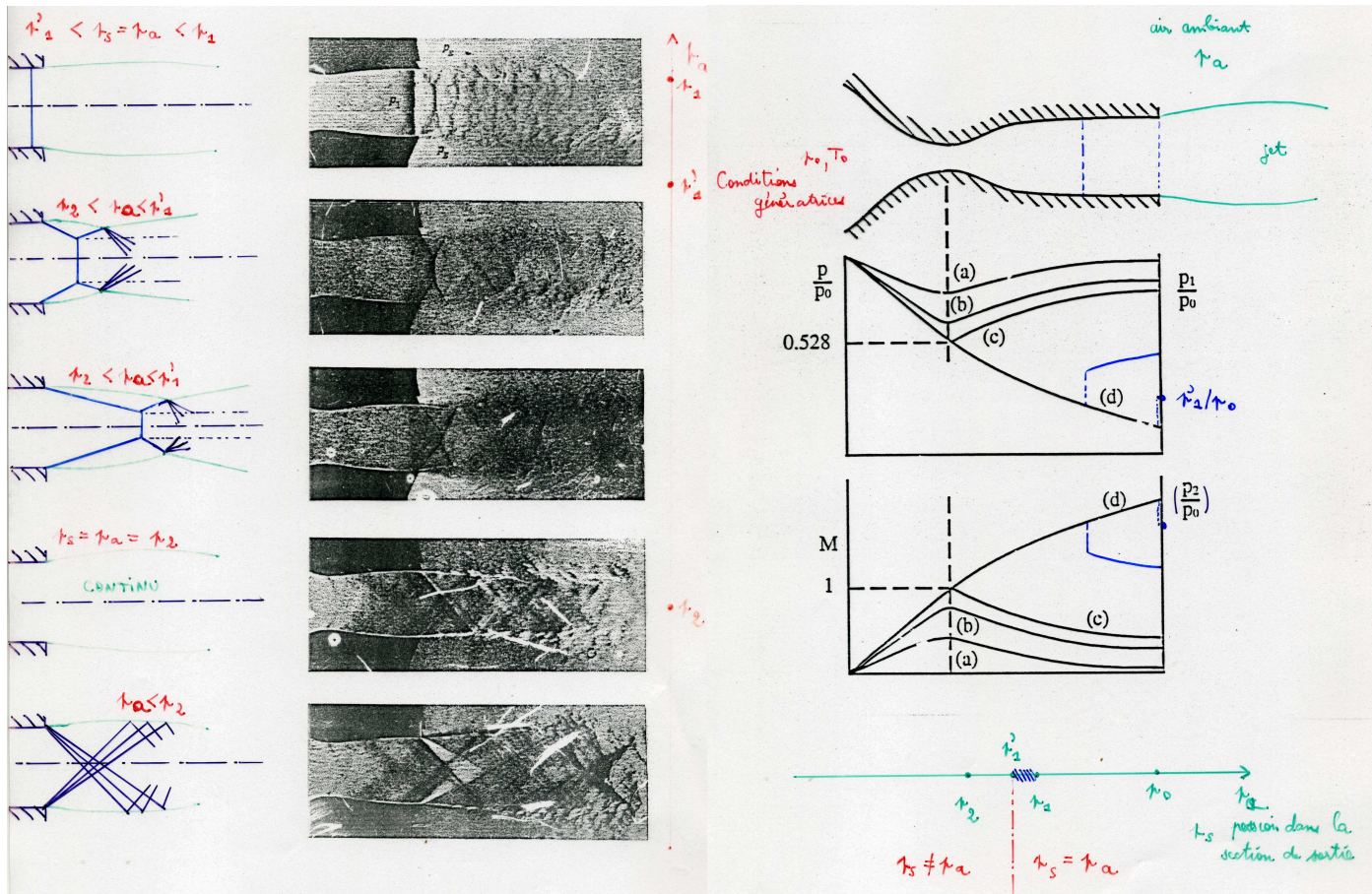


$$\int_{\mathcal{D}'} \{p \underline{n}' + \rho (\underline{u} \cdot \underline{n}') \underline{U}\} da = \underline{0}$$

$$\underline{R} = - \int_{c'} \{p \underline{n} + \rho (\underline{u} \cdot \underline{n}) \underline{u}\} da$$

Requires the far-field behaviour of  $(\rho, p, \underline{u})$

# Momentum conservation for the one-dimensional approximation of a flowing non-viscous fluid



# FLOWING NON-VISCOUS FLUID: BERNOULLI THEOREMS

First case

- potential flow:  $\underline{u} = \underline{grad}[\phi]$  (then  $\underline{rot}(\underline{u}) = \underline{0}$ )
- conservative body force:  $\underline{F} = -\underline{grad}[\Phi]$  (for example for the gravity,  $\Phi = -\underline{g} \cdot \underline{x}$ )
- non-viscous fluid

$$\frac{\partial \underline{u}}{\partial t} + \underline{grad}\left[\frac{u^2}{2}\right] + \underline{rot}(\underline{u}) \wedge \underline{u} = \underline{F} - \left[\frac{1}{\rho}\right] \underline{grad}[p]$$

which, under the assumptions, becomes

$$\underline{grad}\left[\frac{\partial \phi}{\partial t} + \frac{u^2}{2} + \Phi\right] + \left[\frac{1}{\rho}\right] \underline{grad}[p] = \underline{0}$$

- (i) **homogeneous** fluid ( $\rho$  uniform). Then the flow is incompressible so that  $\Delta\phi = 0$  and

$$\underline{grad}\left[\frac{\partial \phi}{\partial t} + \frac{u^2}{2} + \Phi + p/\rho\right] = \underline{0}, \quad \frac{\partial \phi}{\partial t} + \frac{u^2}{2} + \Phi + p/\rho = C(t) \quad \text{Bernoulli theorem}$$

- (ii) **compressible** ( $\rho$  non-uniform) and **barotropic** flow

$$\rho = \rho(p), \quad \frac{\partial \phi}{\partial t} + \frac{u^2}{2} + \Phi + \int dp/\rho = C(t) \quad \text{Bernoulli theorem}$$

# FLOWING NON-VISCOUS FLUID: BERNOULLI THEOREMS

Second case

- steady flow
- conservative body force:  $\underline{F} = -\underline{grad}[\Phi]$
- barotropic flow, i.e.  $\rho = \rho(p)$

Then

- trajectories=streamlines. On such a curve, the elementary displacement  $d\underline{x}$  is parallel with  $\underline{u}$  and

$$\underline{grad}\left[\frac{u^2}{2} + \Phi\right] \cdot d\underline{x} + \left[\frac{1}{\rho}\right] \underline{grad}[p] \cdot d\underline{x} = 0$$

- (i) **homogeneous** fluid ( $\rho$  uniform). Then the flow is incompressible and

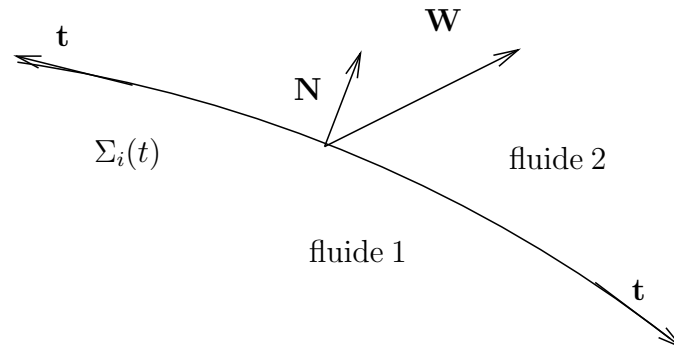
$$\underline{grad}\left[\frac{u^2}{2} + \Phi + p/\rho\right] = \underline{0}, \quad \frac{u^2}{2} + \Phi + p/\rho = C_{streamline}, \quad \text{Bernoulli theorem}$$

- (ii) **compressible** ( $\rho$  non-uniform) and **barotropic** flow

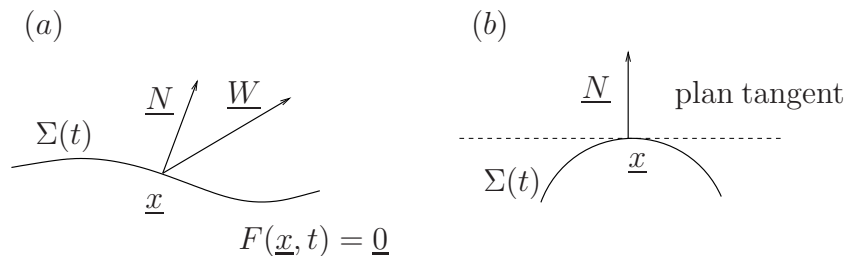
$$\underline{grad}\left[\frac{u^2}{2} + \Phi + \int dp/\rho\right] = \underline{0}, \quad \frac{u^2}{2} + \Phi + \int dp/\rho = C_{streamline}, \quad \text{Bernoulli theorem}$$

Two illustrating applications: the Venturi tube and the Torricelli formula

# Fluid-fluid interface (liquid-gas or liquid-liquid interface)



- immiscible fluids
- $\Sigma_i(t)$  has its own velocity  $\underline{W}$
- Henceforth,  $[[g]] = g(2) - g(1)$  on the interface
- $\underline{N}$  unit normal adequately directed depending upon the interface local concavity and directed from fluid 1 to fluid 2

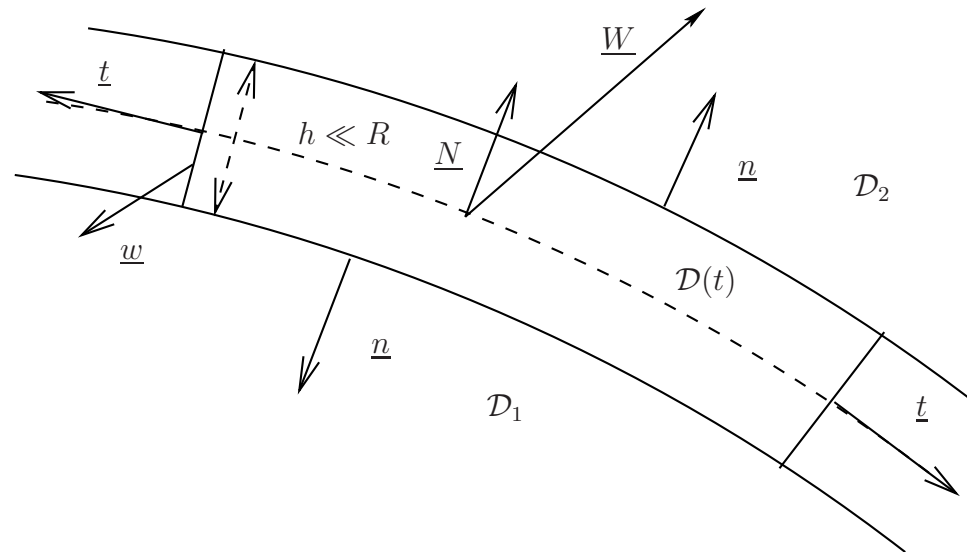


- Local mean curvature  $H$

$$\operatorname{div}_S(\underline{N}) = 2H, \quad 2H = 1/R_1 + 1/R_2$$



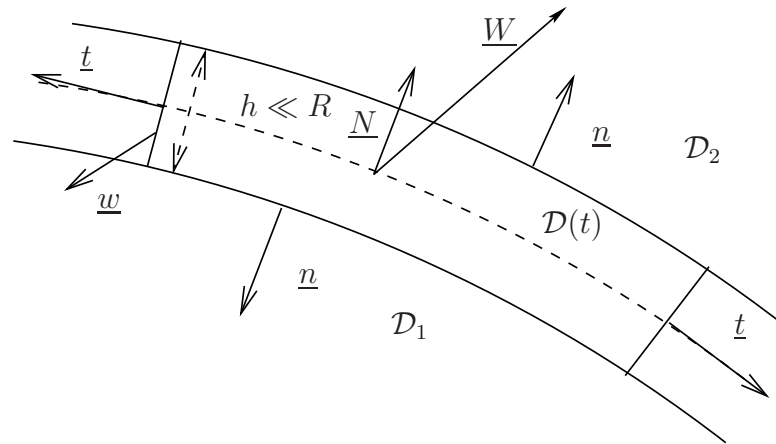
## Mass conservation



- Vanishing thickness  $h \ll R$
- The interface density  $\rho_{int}$  negligible
- Immiscible fluids (gas-liquid, liquid-liquid)
  - Mass conservation

$$\underline{u}_1 \cdot \underline{N} = \underline{u}_2 \cdot \underline{N} = \underline{W} \cdot \underline{N} \quad \text{on } \Sigma_i$$

## Momentum conservation

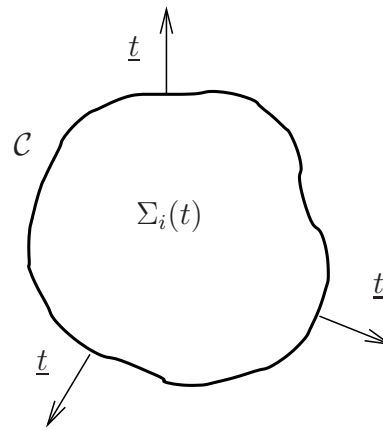


$$\frac{\partial}{\partial t} \int_{\mathcal{D}(t)} \rho_{int} \underline{u} d\Omega - \int_{\mathcal{D}(t)} \rho_{int} \underline{F} d\Omega = \int_{\partial\mathcal{D}_1 \cup \partial\mathcal{D}_2} \underline{\underline{\sigma}} \cdot \underline{n} dS + \int_{\partial\mathcal{D}_l} \underline{\underline{\sigma}} \cdot \underline{n} dS - \int_{\partial\mathcal{D}(t)} [\rho(\underline{u} - \underline{w}) \cdot \underline{n}] \underline{u} dS$$

- The interface density is negligible ( $\rho_{int} \sim 0$ ) so that

$$\int_{\Sigma_i} \underline{\underline{\underline{\sigma}}} \cdot \underline{N} dS + \lim_{h \rightarrow 0} \int_{\partial\mathcal{D}_l} \underline{\underline{\sigma}} \cdot \underline{n} dS = \underline{0}$$

## Surface tension



$$\lim_{h \rightarrow 0} \int_{\partial \mathcal{D}_h} \underline{\underline{\sigma}} \cdot \underline{n} dS = \oint_c \gamma \underline{t} dl$$

- $\gamma \geq 0$  is the interface **surface tension** (given in  $N/m$ )
  - **Auxiliary mathematical identity**

$$\oint_c \gamma \underline{t} dl = \int_{\Sigma_i} \{ \underline{\underline{grad}}_S[\gamma] - \gamma \text{div}_S(\underline{N}) \underline{N} \} dS$$

## Global and local results

$$\int_{\Sigma_i} \{ \llbracket \underline{\sigma} \rrbracket \cdot \underline{N} + \underline{grad}_S[\gamma] - \gamma \operatorname{div}_S(\underline{N}) \underline{N} \} dS = \underline{0}$$

$$\llbracket \underline{\sigma} \rrbracket \cdot \underline{N} + \underline{grad}_S[\gamma] - 2\gamma H \underline{N} = \underline{0} \quad \text{sur } \Sigma_i$$

In summary, at each point on the interface

$$\underline{N} \cdot \llbracket \underline{\sigma} \rrbracket \cdot \underline{N} = 2\gamma H \quad \text{sur } \Sigma_i$$

$$(\llbracket \underline{\sigma} \rrbracket \cdot \underline{N} + \underline{grad}_S[\gamma]) \wedge \underline{N} = \underline{0} \quad \text{on } \Sigma_i$$

$$\underline{u}_1 \cdot \underline{N} = \underline{u}_2 \cdot \underline{N} = \underline{W} \cdot \underline{N} \quad \text{on } \Sigma_i$$

- With  $H \geq 0$ ,  $\underline{N}$  directed from fluid 1 to fluid 2 and  $\llbracket g \rrbracket = g(2) - g(1)$ 
  - Comments

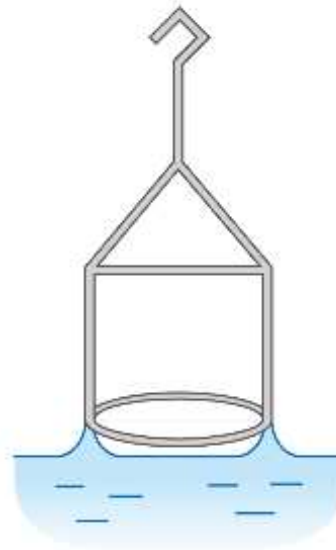
## Experimental values of the surface tension versus the temperature

**Table 1.4.2** Surface Tension as a Function of Temperature for Some Fluid Pairs. (Data from Jasper, J.J. 1972. The Surface Tension of Pure Liquid Compounds. *J. Phys. Chem. Ref. Data* 1, No. 4, 841–1009.)

Fluid Pair	$\sigma$ mN/m	
Acetone/Air	$26.26 - 0.1120T \pm 0.15$	$25 \leq T(C) \leq 50$
Benzaldehyde/Air	$40.72 - 0.1090T \pm 0.5$	$10 \leq T(C) \leq 100$
Carbon tetrachloride/Air	$29.49 - 0.1224T$	$15 \leq T(C) \leq 105$
Cyclohexane/Air	$27.62 - 0.1188T \pm 0.1$	$5 \leq T(C) \leq 70$
Decane/Nitrogen	$25.67 - 0.0920T \pm 0.1$	$10 \leq T(C) \leq 120$
Deuterium oxide/Vapor	$80.62 - 0.2201T \pm 0.1$	$100 \leq T(C) \leq 215$
Dibutyl ether/Air	$24.78 - 0.0934T \pm 0.1$	$15 \leq T(C) \leq 90$
Diethylene glycol/Vapor	$46.97 - 0.0880T \pm 0.14$	$20 \leq T(C) \leq 140$
Dodecane/Nitrogen	$27.12 - 0.0884T \pm 0.1$	$10 \leq T(C) \leq 120$
Ethanol/Air	$24.05 - 0.0832T \pm 0.1$	$10 \leq T(C) \leq 70$
Ethyl acetate/Air	$26.29 - 0.1161T \pm 0.1$	$10 \leq T(C) \leq 100$
Ethylene glycol/Vapor	$50.21 - 0.0890T \pm 0.14$	$20 \leq T(C) \leq 140$
Formamide/Air	$59.13 - 0.0842T \pm 0.2$	$25 \leq T(C) \leq 120$
Heavy water/Air	$74.64 - 0.1082T^{1.1} \pm 0.18$	$10 \leq T(C) \leq 75$
Helium II/Vapor	$0.352 - 0.0069T^{\frac{1}{2}}$	$0.5 \leq T(K) \leq 2.1$
Hexane/Nitrogen	$20.44 - 0.1020T \pm 0.1$	$10 \leq T(C) \leq 60$
Isopropyl alcohol/Air	$22.90 - 0.0789T \pm 0.1$	$10 \leq T(C) \leq 80$
Mercury/Vapor	$490.6 - 0.2049T \pm 2.0$	$5 \leq T(C) \leq 200$
Methanol/Air	$24.00 - 0.0773T \pm 0.1$	$10 \leq T(C) \leq 60$
Octane/Nitrogen	$23.52 - 0.0951T \pm 0.1$	$10 \leq T(C) \leq 120$
Toluene/Air	$30.90 - 0.1189T \pm 0.1$	$10 \leq T(C) \leq 100$
Water/Air	$75.83 - 0.1477T \pm 0.1$	$10 \leq T(C) \leq 100$
Xenon/Vapor	$0.03703(289.74 - T)^{1.287}$	$165 \leq T(K) \leq 285$

C = degrees Celsius; K = Kelvin.

How to experimentally measure  $\gamma$ ?  
Examples: Jurin's law, Wilhelmy's ring



$$\gamma = \frac{F}{4\pi R} C$$

Where:

$\gamma$  : Surface tension

F : Force acting on the ring

R : Center diameter of the ring

C : Correction factor